

Heat and Mass Transfer, 3rd Edition

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Chapter 8

INTERNAL FORCED CONVECTION

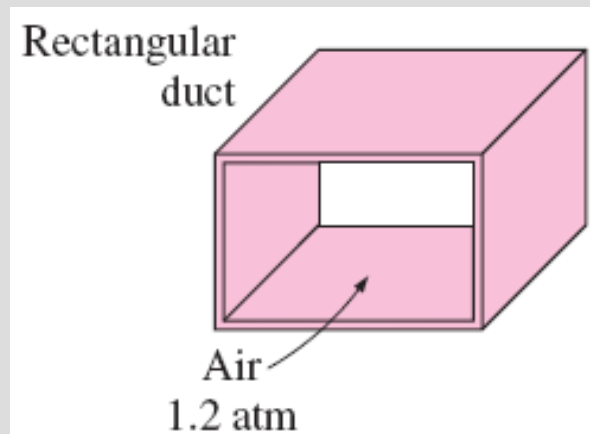
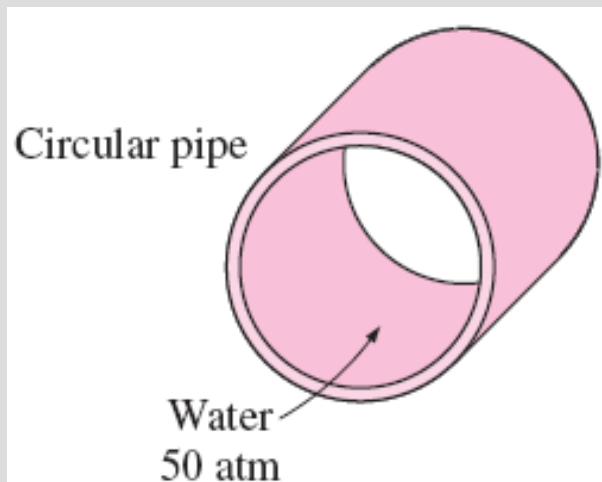
Mehmet Kanoglu

Objectives

- Obtain average velocity from a knowledge of velocity profile, and average temperature from a knowledge of temperature profile in internal flow,
- Have a visual understanding of different flow regions in internal flow, and calculate hydrodynamic and thermal entry lengths
- Analyze heating and cooling of a fluid flowing in a tube under constant surface temperature and constant surface heat flux conditions, and work with the logarithmic mean temperature difference
- Obtain analytic relations for the velocity profile, pressure drop, friction factor, and Nusselt number in fully developed laminar flow, and
- Determine the friction factor and Nusselt number in fully developed turbulent flow using empirical relations, and calculate the heat transfer rate.

INTRODUCTION

- Liquid or gas flow through *pipes* or *ducts* is commonly used in heating and cooling applications and fluid distribution networks.
- The fluid in such applications is usually forced to flow by a fan or pump through a flow section.
- We pay particular attention to *friction*, which is directly related to the *pressure drop* and *head loss* during flow through pipes and ducts.
- The pressure drop is then used to determine the *pumping power requirement*.



Circular pipes can withstand large pressure differences between the inside and the outside without undergoing any significant distortion, but noncircular pipes cannot.

Theoretical solutions are obtained only for a few simple cases such as fully developed laminar flow in a circular pipe.

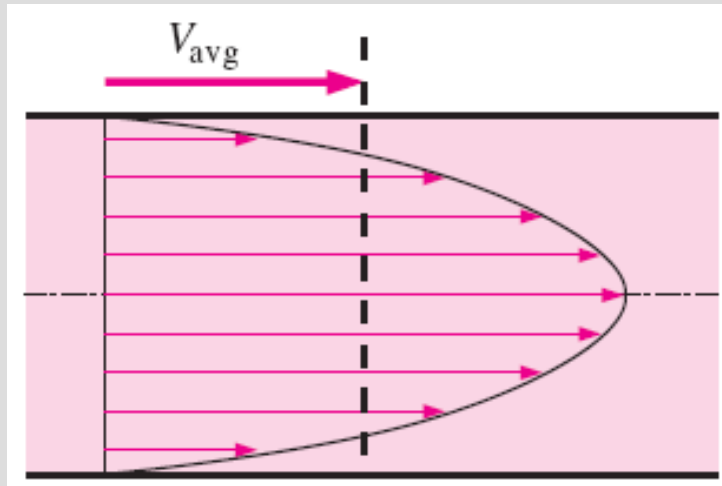
Therefore, we must rely on experimental results and empirical relations for most fluid flow problems rather than closed-form analytical solutions.

$$\dot{m} = \rho V_{\text{avg}} A_c = \int_{A_c} \rho u(r) dA_c$$

The value of the average velocity V_{avg} at some streamwise cross-section is determined from the requirement that the *conservation of mass* principle be satisfied

$$V_{\text{avg}} = \frac{\int_{A_c} \rho u(r) dA_c}{\rho A_c} = \frac{\int_0^R \rho u(r) 2\pi r dr}{\rho \pi R^2} = \frac{2}{R^2} \int_0^R u(r) r dr$$

The average velocity for incompressible flow in a circular pipe of radius R



Average velocity V_{avg} is defined as the average speed through a cross section. For fully developed laminar pipe flow, V_{avg} is half of the maximum velocity.

GENERAL CONSIDERATIONS FOR PIPE FLOW

Liquid or gas flow through pipes or ducts is commonly used in practice in heating and cooling applications. The fluid is forced to flow by a fan or pump through a conduit that is sufficiently long to accomplish the desired heat transfer.

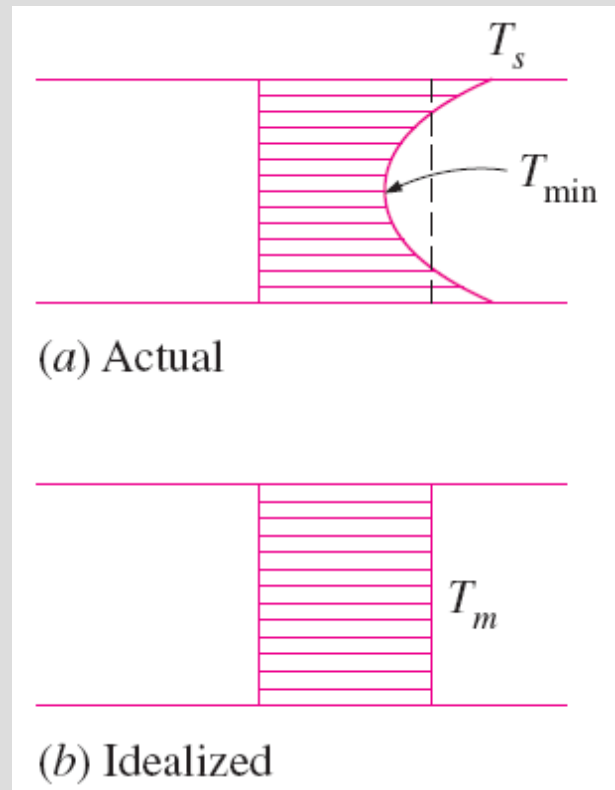
Transition from laminar to turbulent flow depends on the Reynolds number as well as the degree of disturbance of the flow by *surface roughness, pipe vibrations, and the fluctuations in the flow.*

The flow in a pipe is laminar for $Re < 2300$, fully turbulent for $Re > 10,000$, and transitional in between.

$$Re = \frac{\rho V_{avg} D}{\mu} = \frac{V_{avg} D}{\nu}$$

$$\dot{E}_{fluid} = \dot{m} c_p T_m = \int_{\dot{m}} c_p T(r) \delta \dot{m} = \int_{A_c} \rho c_p T(r) u(r) V dA_c$$

$$T_m = \frac{\int_{\dot{m}} c_p T(r) \delta \dot{m}}{\dot{m} c_p} = \frac{\int_0^R c_p T(r) \rho u(r) 2\pi r dr}{\rho V_{avg} (\pi R^2) c_p} = \frac{2}{V_{avg} R^2} \int_0^R T(r) u(r) r dr$$



Actual and idealized temperature profiles for flow in a tube (the rate at which energy is transported with the fluid is the same for both cases).

For flow through noncircular pipes, the Reynolds number is based on the **hydraulic diameter**

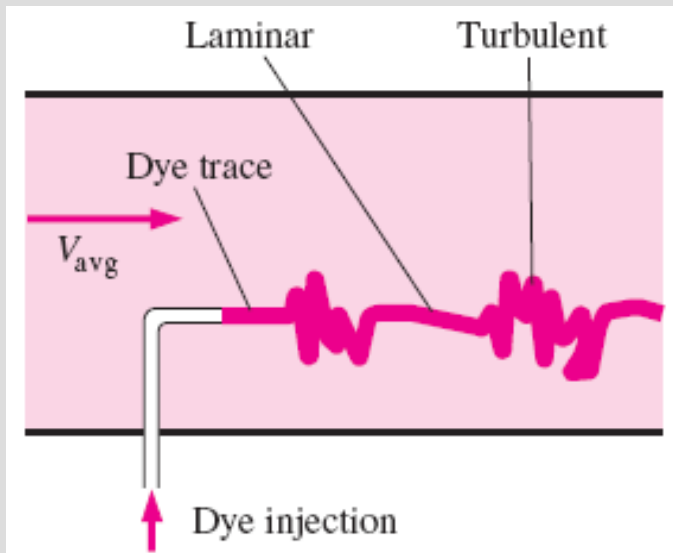
$$D_h = \frac{4A_c}{p}$$

For flow in a circular pipe:

$Re \lesssim 2300$ laminar flow

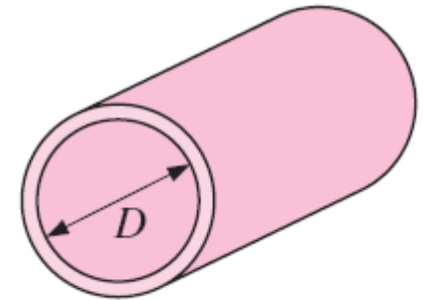
$2300 \lesssim Re \lesssim 10,000$ transitional flow

$Re \gtrsim 10,000$ turbulent flow



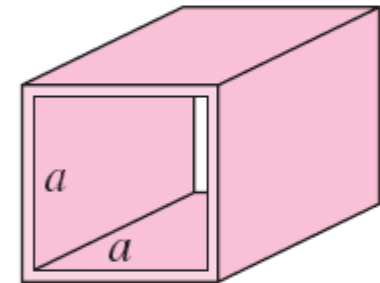
In the transitional flow region of $2300 \leq Re \leq 10,000$, the flow switches between laminar and turbulent seemingly randomly.

Circular tube:



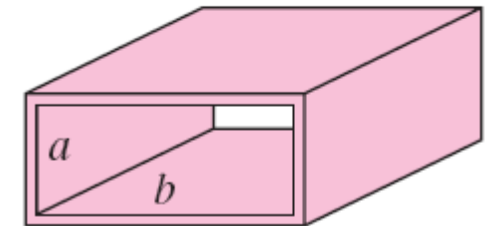
$$D_h = \frac{4(\pi D^2/4)}{\pi D} = D$$

Square duct:



$$D_h = \frac{4a^2}{4a} = a$$

Rectangular duct:



$$D_h = \frac{4ab}{2(a+b)} = \frac{2ab}{a+b}$$

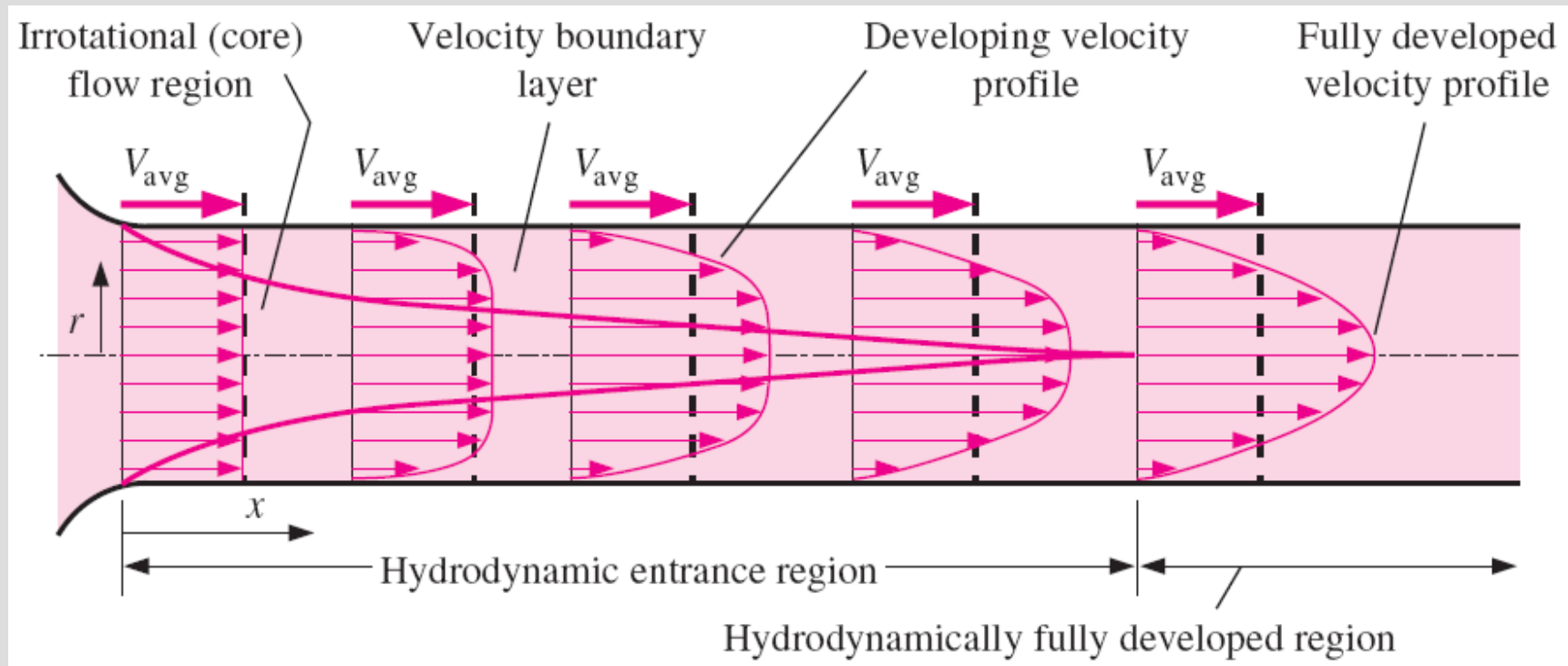
The hydraulic diameter $D_h = 4A_c/p$ is defined such that it reduces to ordinary diameter for circular tubes.

THE ENTRANCE REGION

Velocity boundary layer: The region of the flow in which the effects of the viscous shearing forces caused by fluid viscosity are felt.

Boundary layer region: The viscous effects and the velocity changes are significant.

Irrotational (core) flow region: The frictional effects are negligible and the velocity remains essentially constant in the radial direction.



The development of the velocity boundary layer in a pipe. The developed average velocity profile is parabolic in laminar flow, but somewhat flatter or fuller in turbulent flow.

The fluid properties in internal flow are usually evaluated at the *bulk mean fluid temperature*, which is the arithmetic average of the mean temperatures at the inlet and the exit: $T_b = (T_{m,i} + T_{m,e})/2$

Thermal Entrance Region

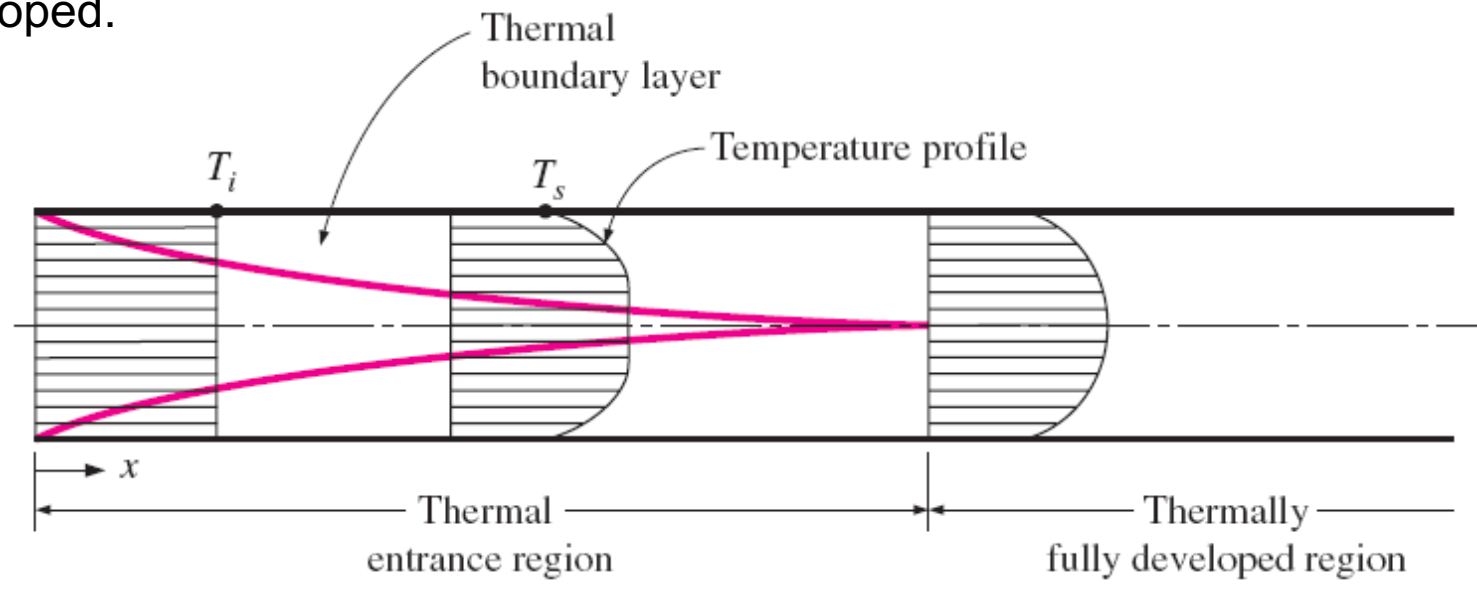
Thermal entrance region: The region of flow over which the thermal boundary layer develops and reaches the tube center.

Thermal entry length: The length of this region.

Thermally developing flow: Flow in the thermal entrance region. This is the region where the temperature profile develops.

Thermally fully developed region: The region beyond the thermal entrance region in which the dimensionless temperature profile remains unchanged.

Fully developed flow: The region in which the flow is both hydrodynamically and thermally developed.



The development of the thermal boundary layer in a tube.

Hydrodynamically fully developed:

$$\frac{\partial u(r, x)}{\partial x} = 0 \quad \longrightarrow \quad u = u(r)$$

Thermally fully developed:

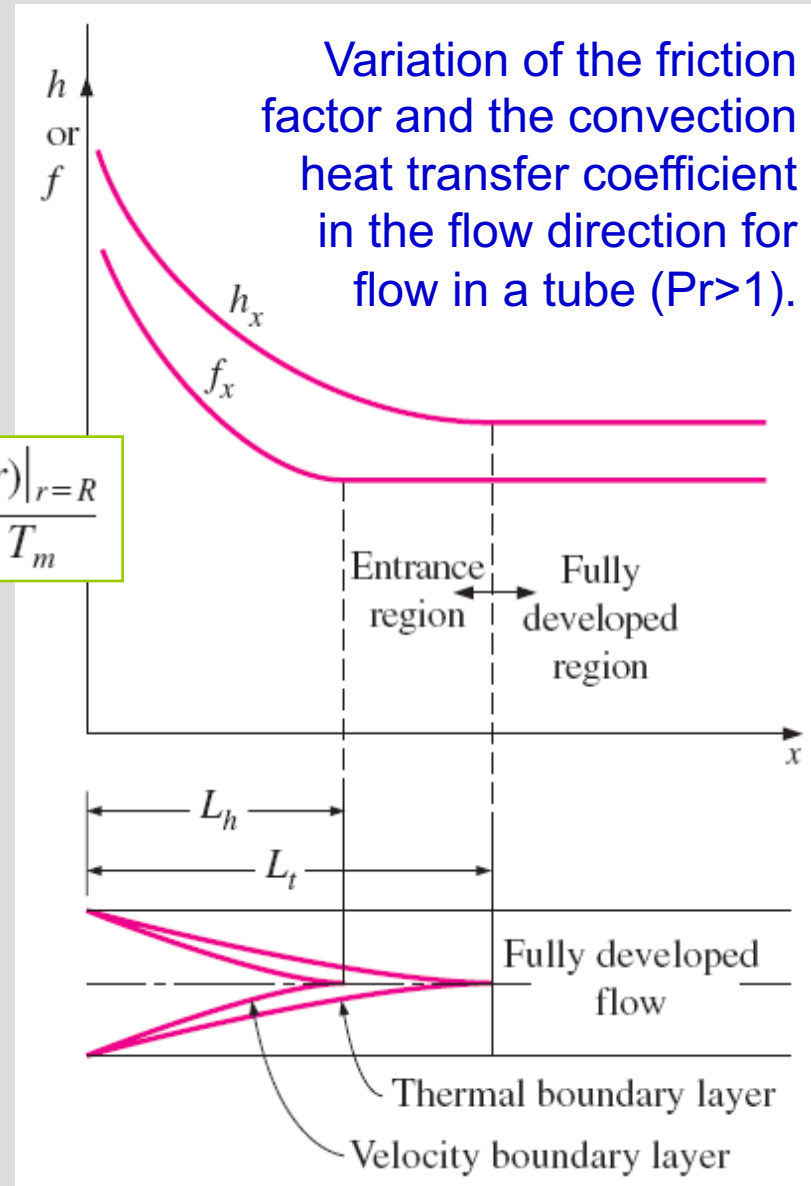
$$\frac{\partial}{\partial x} \left[\frac{T_s(x) - T(r, x)}{T_s(x) - T_m(x)} \right] = 0$$

$$\dot{q}_s = h_x(T_s - T_m) = k \left. \frac{\partial T}{\partial r} \right|_{r=R} \quad \longrightarrow \quad h_x = \frac{k(\partial T / \partial r)|_{r=R}}{T_s - T_m}$$

In the thermally fully developed region of a tube, the local convection coefficient is constant (does not vary with x).

Therefore, both the friction (which is related to wall shear stress) and convection coefficients remain constant in the fully developed region of a tube.

The pressure drop and heat flux are *higher* in the entrance regions of a tube, and the effect of the entrance region is always to *increase* the average friction factor and heat transfer coefficient for the entire tube.



Entry Lengths

$$L_{h, \text{ laminar}} \approx 0.05 \text{ Re } D$$

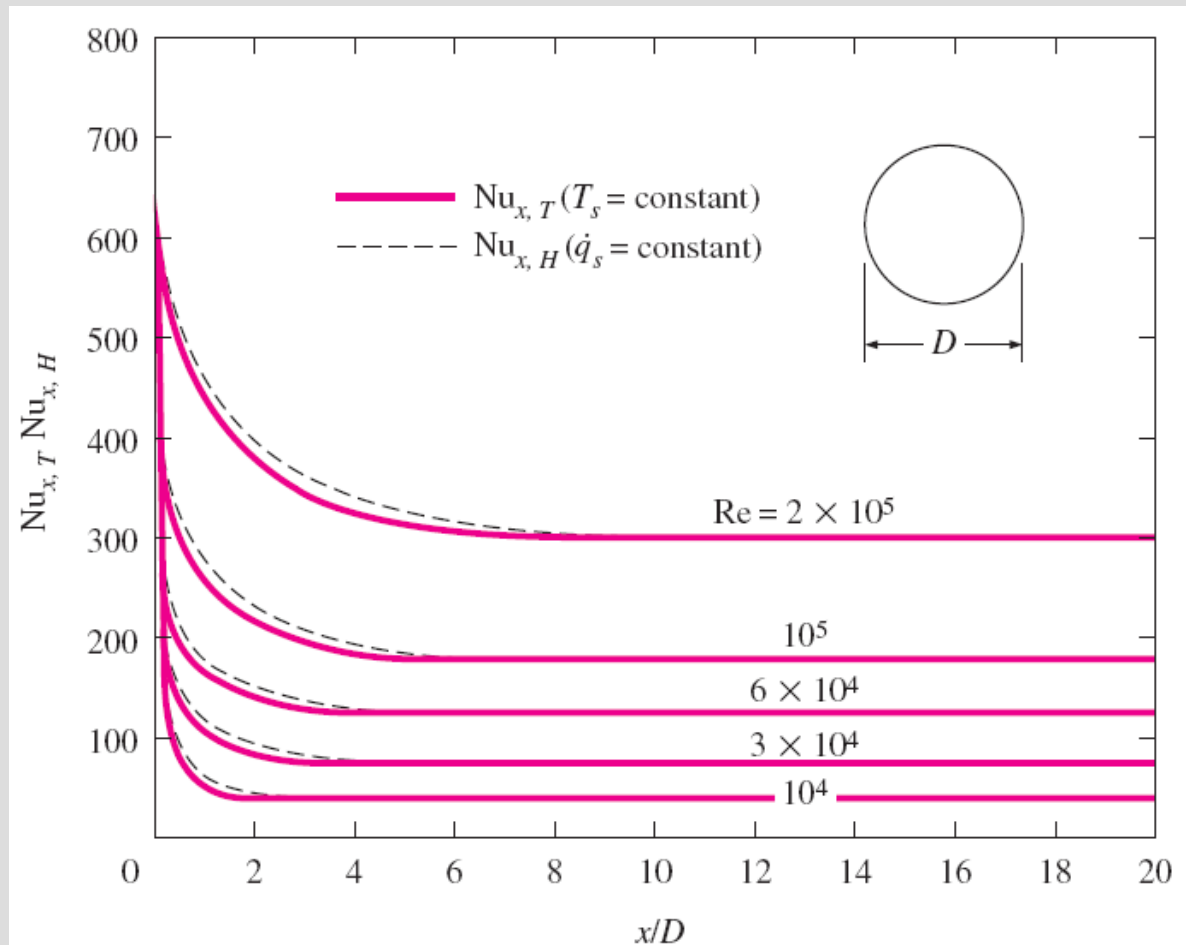
$$L_{t, \text{ laminar}} \approx 0.05 \text{ Re Pr } D = \text{Pr } L_{h, \text{ laminar}}$$

$$L_{h, \text{ turbulent}} = 1.359 D \text{ Re}^{1/4}$$

$$L_{h, \text{ turbulent}} \approx L_{t, \text{ turbulent}} \approx 10 D$$

- The Nusselt numbers and thus h values are much higher in the entrance region.
- The Nusselt number reaches a constant value at a distance of less than 10 diameters, and thus the flow can be assumed to be fully developed for $x > 10D$.
- The Nusselt numbers for the uniform surface temperature and uniform surface heat flux conditions are identical in the fully developed regions, and nearly identical in the entrance regions.

Variation of local Nusselt number along a tube in turbulent flow for both uniform surface temperature and uniform surface heat flux.



Entry Lengths

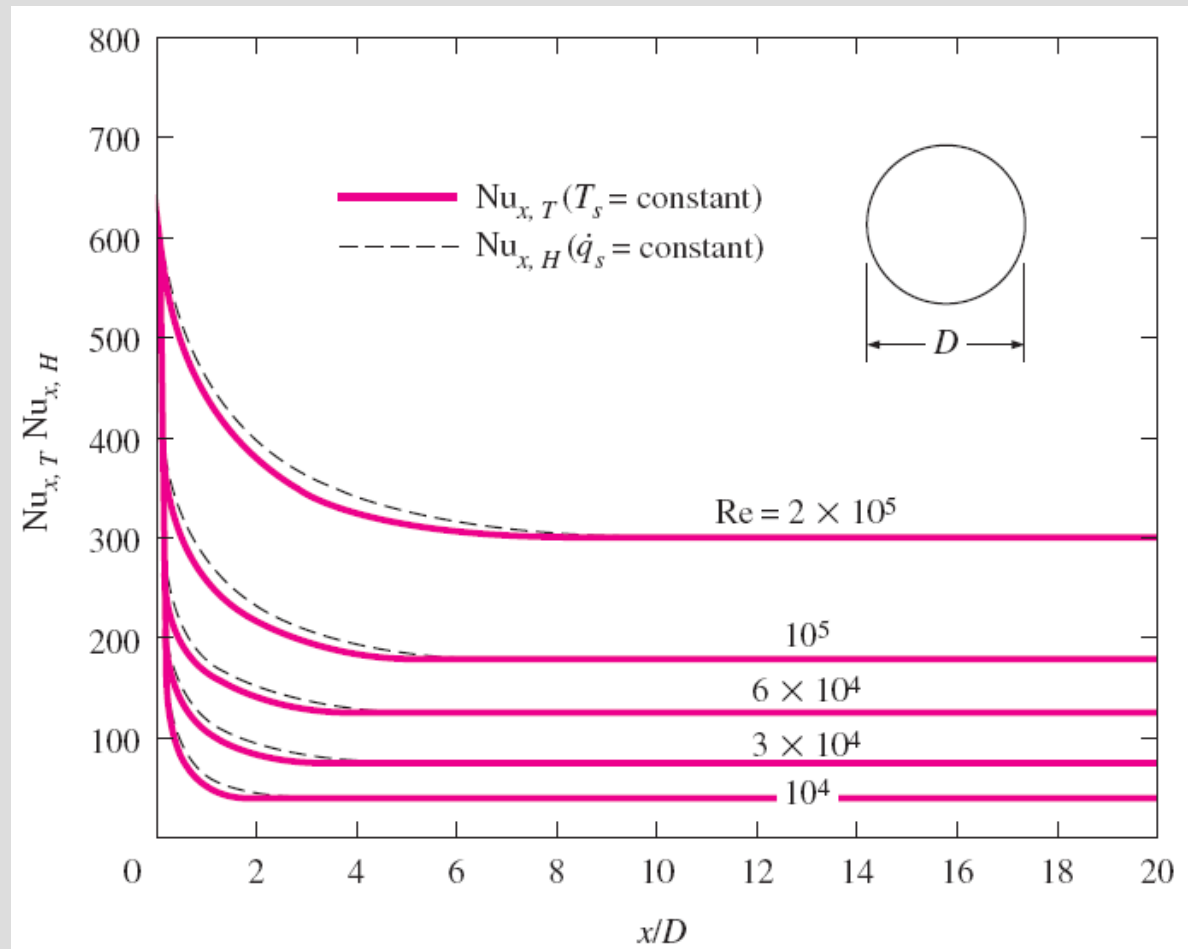
$$L_{h, \text{ laminar}} \approx 0.05 \text{ Re } D$$

$$L_{t, \text{ laminar}} \approx 0.05 \text{ Re Pr } D = \text{Pr } L_{h, \text{ laminar}}$$

$$L_{h, \text{ turbulent}} = 1.359 D \text{ Re}^{1/4}$$

$$L_{h, \text{ turbulent}} \approx L_{t, \text{ turbulent}} \approx 10D$$

Variation of local Nusselt number along a tube in turbulent flow for both uniform surface temperature and uniform surface heat flux.



GENERAL THERMAL ANALYSIS

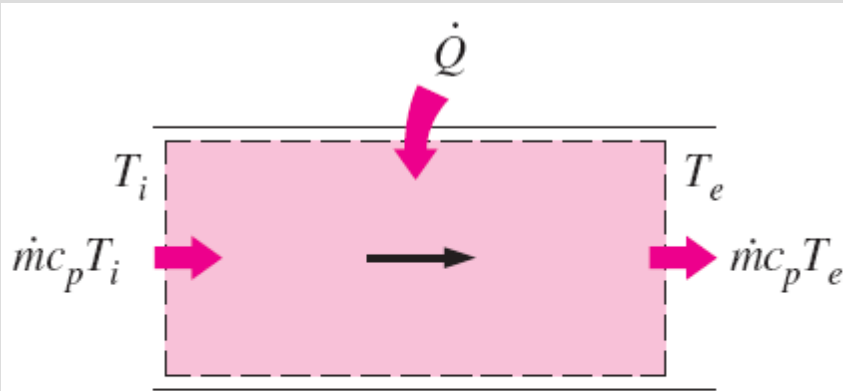
Rate of heat transfer

$$\dot{Q} = \dot{m}c_p(T_e - T_i) \quad (\text{W})$$

Surface heat flux

$$\dot{q}_s = h_x(T_s - T_m) \quad (\text{W/m}^2)$$

h_x the *local* heat transfer coefficient



Energy balance:

$$\dot{Q} = \dot{m}c_p(T_e - T_i)$$

The heat transfer to a fluid flowing in a tube is equal to the increase in the energy of the fluid.

The thermal conditions at the surface can be approximated to be

constant surface temperature ($T_s = \text{const}$)

constant surface heat flux ($q_s = \text{const}$)

The constant surface temperature condition is realized when a phase change process such as boiling or condensation occurs at the outer surface of a tube.

The constant surface heat flux condition is realized when the tube is subjected to radiation or electric resistance heating uniformly from all directions.

We may have either $T_s = \text{constant}$ or $q_s = \text{constant}$ at the surface of a tube, but not both.

Constant Surface Heat Flux ($q_s = \text{constant}$)

Rate of heat transfer:

$$\dot{Q} = \dot{q}_s A_s = \dot{m} c_p (T_e - T_i) \quad (\text{W})$$

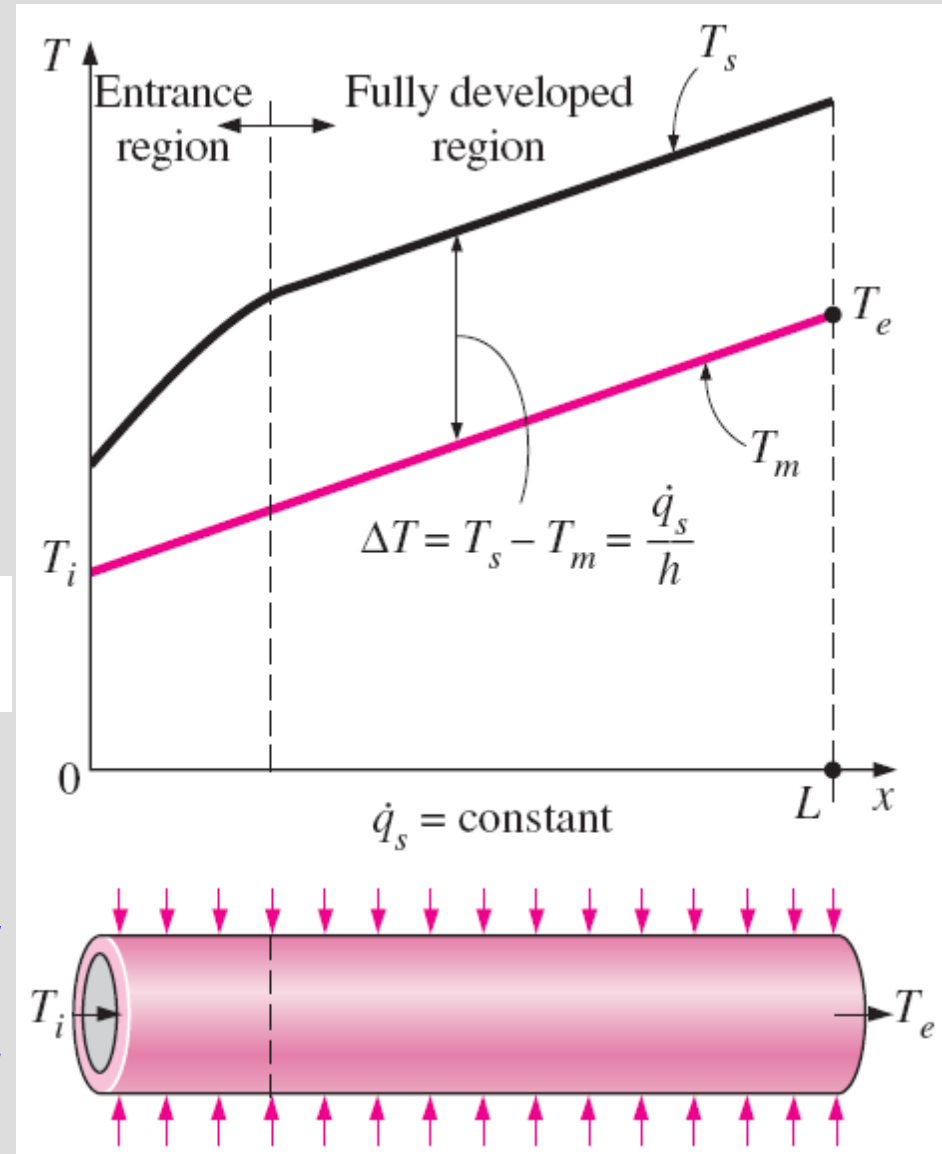
Mean fluid temperature
at the tube exit:

$$T_e = T_i + \frac{\dot{q}_s A_s}{\dot{m} c_p}$$

Surface temperature:

$$\dot{q}_s = h(T_s - T_m) \longrightarrow T_s = T_m + \frac{\dot{q}_s}{h}$$

Variation of the *tube surface* and the *mean fluid* temperatures along the tube for the case of constant surface heat flux.



$$\dot{m} c_p dT_m = \dot{q}_s(p dx) \longrightarrow \frac{dT_m}{dx} = \frac{\dot{q}_s p}{\dot{m} c_p} = \text{constant}$$

$$\frac{dT_m}{dx} = \frac{dT_s}{dx}$$

$$\frac{\partial}{\partial x} \left(\frac{T_s - T}{T_s - T_m} \right) = 0 \longrightarrow \frac{1}{T_s - T_m} \left(\frac{\partial T_s}{\partial x} - \frac{\partial T}{\partial x} \right) = 0$$

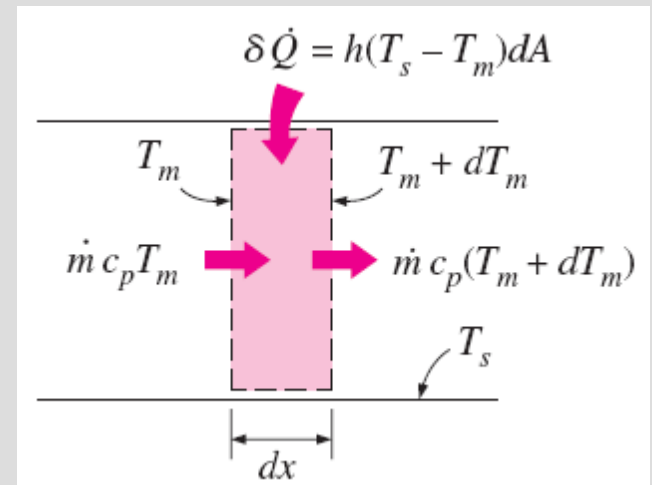
$$\longrightarrow \frac{\partial T}{\partial x} = \frac{dT_s}{dx}$$

$$\frac{\partial T}{\partial x} = \frac{dT_s}{dx} = \frac{dT_m}{dx} = \frac{\dot{q}_s p}{\dot{m} c_p} = \text{constant}$$

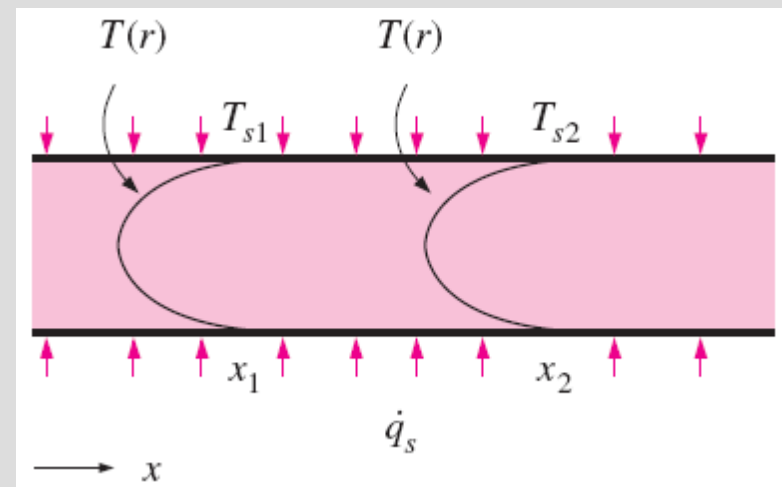
Circular tube:

$$\frac{\partial T}{\partial x} = \frac{dT_s}{dx} = \frac{dT_m}{dx} = \frac{2\dot{q}_s}{\rho V_{\text{avg}} c_p R} = \text{constant}$$

The shape of the temperature profile remains unchanged in the fully developed region of a tube subjected to constant surface heat flux.



Energy interactions for a differential control volume in a tube.



Constant Surface Temperature ($T_s = \text{constant}$)

Rate of heat transfer to or from a fluid flowing in a tube

$$\dot{Q} = hA_s \Delta T_{\text{avg}} = hA_s (T_s - T_m)_{\text{avg}} \quad (\text{W})$$

Two suitable ways of expressing ΔT_{avg}

- arithmetic mean temperature difference
- logarithmic mean temperature difference

Arithmetic mean temperature difference

$$\Delta T_{\text{avg}} \approx \Delta T_{\text{am}} = \frac{\Delta T_i + \Delta T_e}{2} = \frac{(T_s - T_i) + (T_s - T_e)}{2} = T_s - \frac{T_i + T_e}{2} = T_s - T_b$$

Bulk mean fluid temperature: $T_b = (T_i + T_e)/2$

By using arithmetic mean temperature difference, we assume that the mean fluid temperature varies linearly along the tube, which is hardly ever the case when $T_s = \text{constant}$.

This simple approximation often gives acceptable results, but not always.

Therefore, we need a better way to evaluate ΔT_{avg} .

$$\dot{m}c_p dT_m = h(T_s - T_m)dA_s$$

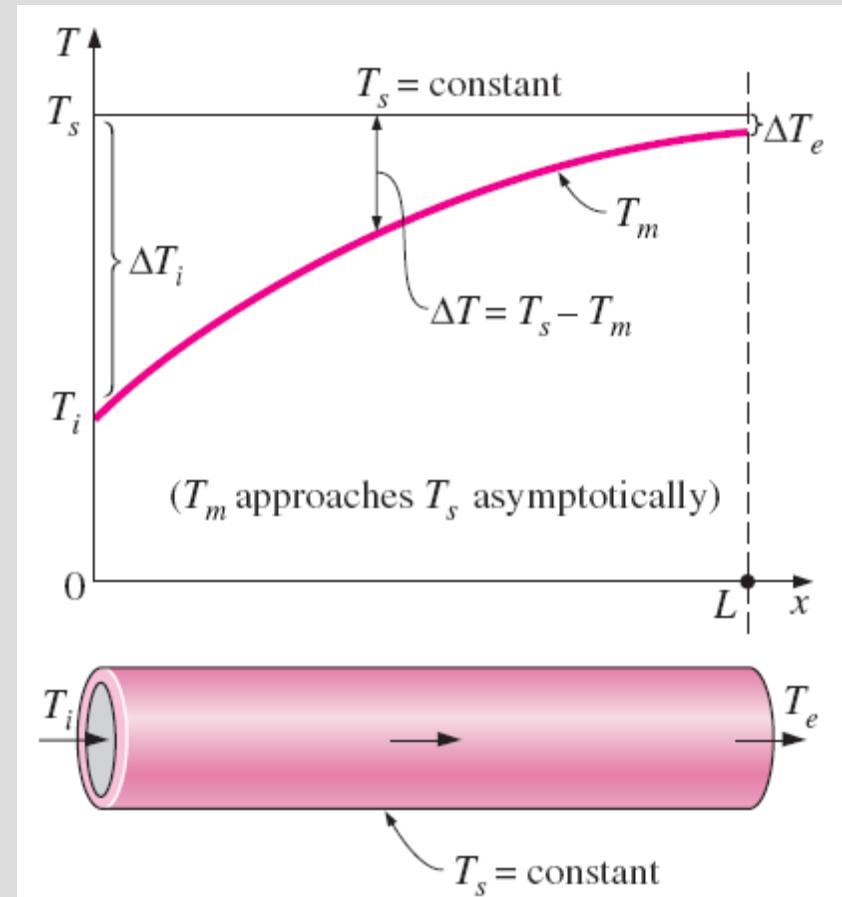
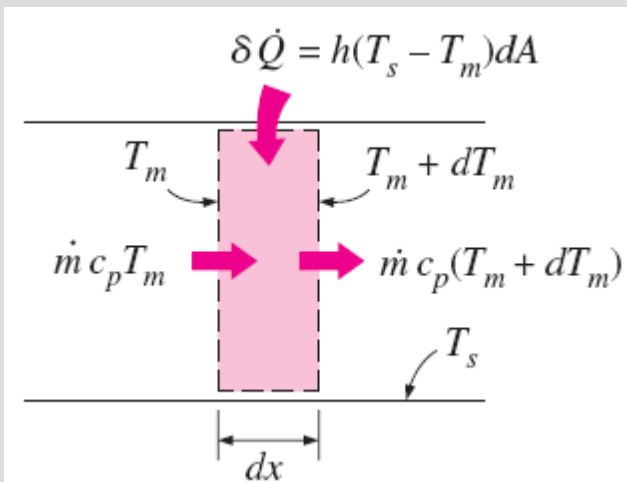
$$dA_s = p dx \quad dT_m = -d(T_s - T_m)$$

$$\frac{d(T_s - T_m)}{T_s - T_m} = -\frac{hp}{\dot{m}c_p} dx$$

Integrating from $x = 0$ (tube inlet, $T_m = T_i$) to $x = L$ (tube exit, $T_m = T_e$)

$$\ln \frac{T_s - T_e}{T_s - T_i} = -\frac{hA_s}{\dot{m}c_p}$$

$$T_e = T_s - (T_s - T_i) \exp(-hA_s/\dot{m}c_p)$$



The variation of the *mean fluid* temperature along the tube for the case of constant temperature.

Energy interactions for a differential control volume in a tube.

$$\dot{Q} = hA_s\Delta T_{\ln}$$

$$\Delta T_{\ln} = \frac{T_i - T_e}{\ln[(T_s - T_e)/(T_s - T_i)]} = \frac{\Delta T_e - \Delta T_i}{\ln(\Delta T_e/\Delta T_i)}$$

logarithmic
mean
temperature
difference

NTU: Number of transfer units. A measure of the effectiveness of the heat transfer systems.

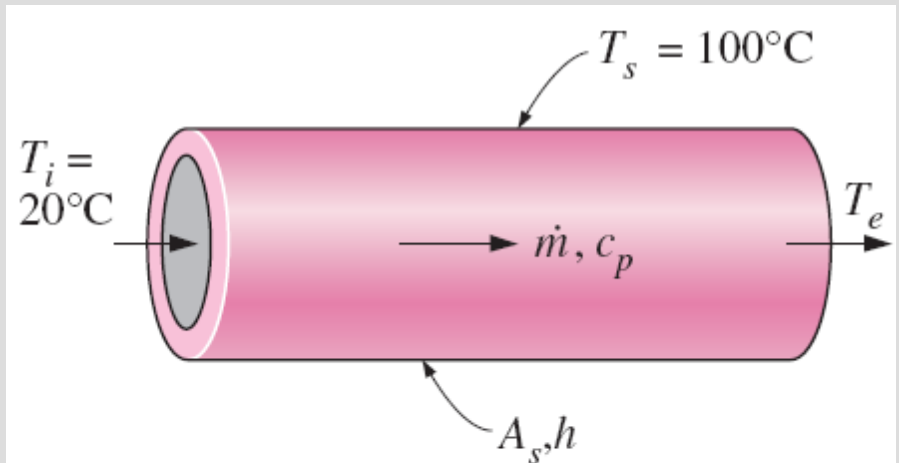
For $NTU = 5$, $T_e = T_s$, and the limit for heat transfer is reached.

A small value of NTU indicates more opportunities for heat transfer.

ΔT_{\ln} is an *exact* representation of the *average temperature difference* between the fluid and the surface.

When ΔT_e differs from ΔT_i by no more than 40 percent, the error in using the arithmetic mean temperature difference is less than 1 percent.

An NTU greater than 5 indicates that the fluid flowing in a tube will reach the surface temperature at the exit regardless of the inlet temperature.



$NTU = hA_s / \dot{m}c_p$	$T_e, ^\circ\text{C}$
0.01	20.8
0.05	23.9
0.10	27.6
0.50	51.5
1.00	70.6
5.00	99.5
10.00	100.0

LAMINAR FLOW IN TUBES

$$\dot{m}c_p T_x - \dot{m}c_p T_{x+dx} + \dot{Q}_r - \dot{Q}_{r+dr} = 0$$

$$\dot{m} = \rho u A_c = \rho u (2\pi r dr)$$

$$\rho c_p u \frac{T_{x+dx} - T_x}{dx} = -\frac{1}{2\pi r dx} \frac{\dot{Q}_{r+dr} - \dot{Q}_r}{dr}$$

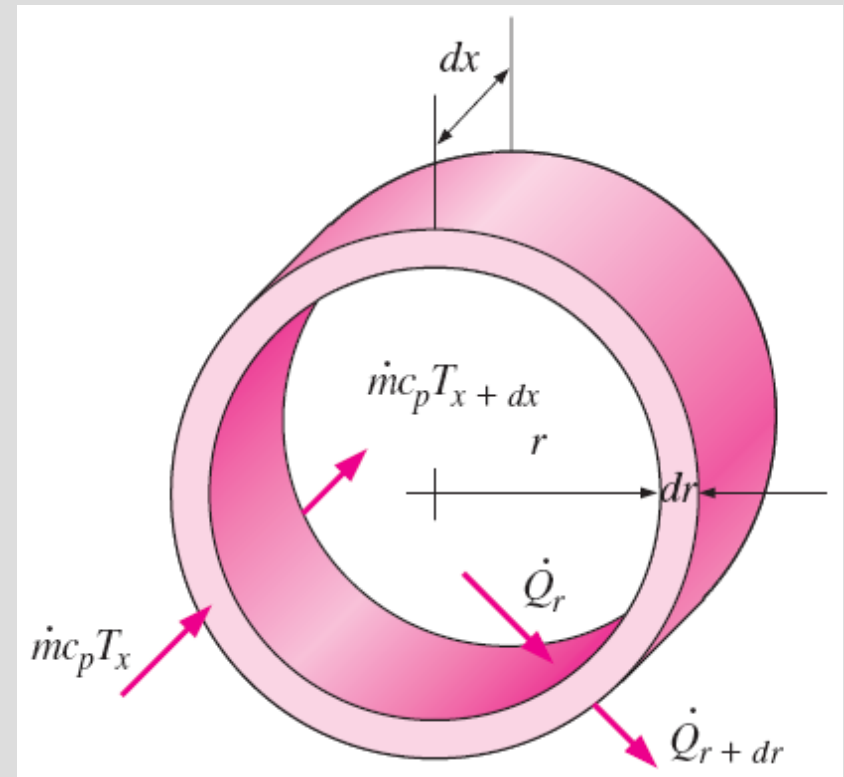
$$u \frac{\partial T}{\partial x} = -\frac{1}{2\rho c_p \pi r dx} \frac{\partial \dot{Q}}{\partial r}$$

$$\frac{\partial \dot{Q}}{\partial r} = \frac{\partial}{\partial r} \left(-k 2\pi r dx \frac{\partial T}{\partial r} \right) = -2\pi k dx \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right)$$

$$\alpha = k/\rho c_p$$

$$u \frac{\partial T}{\partial x} = \frac{\alpha}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right)$$

The rate of net energy transfer to the control volume by mass flow is equal to the net rate of heat conduction in the radial direction.



The differential volume element used in the derivation of energy balance relation.

Constant Surface Heat Flux

$$\frac{\partial T}{\partial x} = \frac{dT_s}{dx} = \frac{dT_m}{dx} = \frac{2\dot{q}_s}{\rho V_{\text{avg}} c_p R} = \text{constant}$$

$$\frac{4\dot{q}_s}{kR} \left(1 - \frac{r^2}{R^2}\right) = \frac{1}{r} \frac{d}{dr} \left(r \frac{dT}{dr} \right)$$

$$T = \frac{\dot{q}_s}{kR} \left(r^2 - \frac{r^4}{4R^2} \right) + C_1 r + C_2$$

Applying the boundary conditions $\partial T / \partial r = 0$ at $r = 0$ (because of symmetry) and $T = T_s$ at $r = R$

$$T = T_s - \frac{\dot{q}_s R}{k} \left(\frac{3}{4} - \frac{r^2}{R^2} + \frac{r^4}{4R^4} \right)$$

$$T_m = T_s - \frac{11}{24} \frac{\dot{q}_s R}{k}$$

$$\dot{q}_s = h(T_s - T_m)$$

$$h = \frac{24}{11} \frac{k}{R} = \frac{48}{11} \frac{k}{D} = 4.36 \frac{k}{D}$$

Circular tube, laminar ($\dot{q}_s = \text{constant}$):

$$\text{Nu} = \frac{hD}{k} = 4.36$$

Therefore, for fully developed laminar flow in a circular tube subjected to constant surface heat flux, the Nusselt number is a constant.

There is no dependence on the Reynolds or the Prandtl numbers.

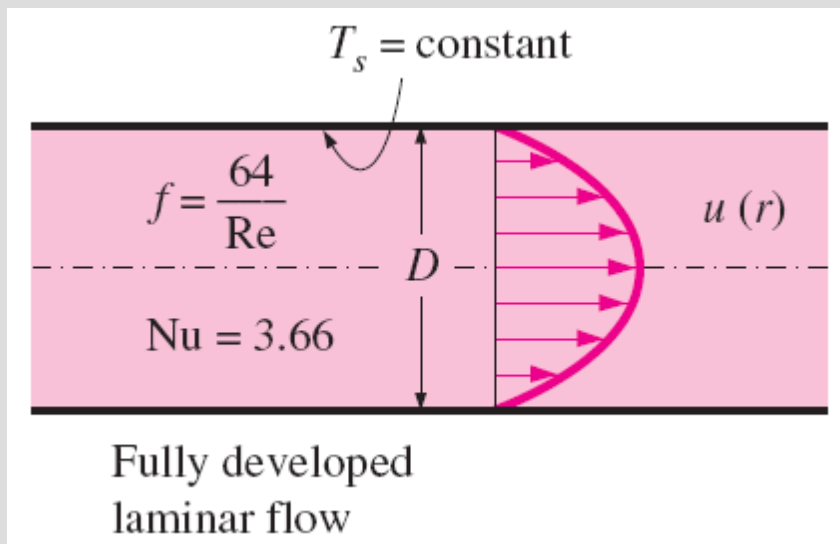
Constant Surface Temperature

Circular tube, laminar ($T_s = \text{constant}$):

$$\text{Nu} = \frac{hD}{k} = 3.66$$

The thermal conductivity k for use in the Nu relations should be evaluated at the bulk mean fluid temperature.

For laminar flow, the effect of *surface roughness* on the friction factor and the heat transfer coefficient is negligible.



In laminar flow in a tube with constant surface temperature, both the *friction factor* and the *heat transfer coefficient* remain constant in the fully developed region.

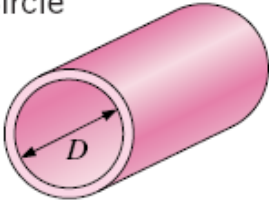
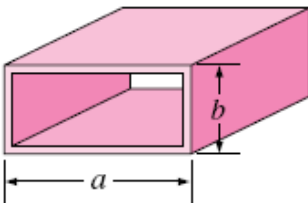
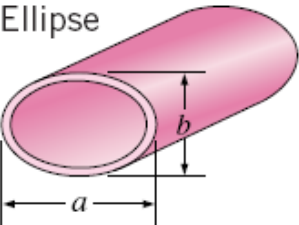
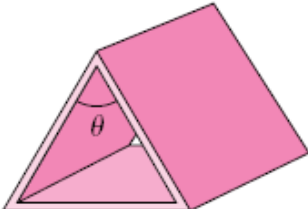
Laminar Flow in Noncircular Tubes

Nusselt number relations are given in the table for *fully developed laminar flow* in tubes of various cross sections.

The Reynolds and Nusselt numbers for flow in these tubes are based on the *hydraulic diameter* $D_h = 4A_c/p$,

Once the Nusselt number is available, the convection heat transfer coefficient is determined from $h = k\text{Nu}/D_h$.

Nusselt number and friction factor for fully developed laminar flow in tubes of various cross sections ($D_h = 4A_c/p$, $Re = V_{avg}D_h/\nu$, and $Nu = hD_h/k$)

Tube Geometry	a/b or θ°	Nusselt Number	
		$T_s = \text{Const.}$	$\dot{q}_s = \text{Const.}$
Circle 	—	3.66	4.36
Rectangle 	a/b		
	1	2.98	3.61
	2	3.39	4.12
	3	3.96	4.79
	4	4.44	5.33
	6	5.14	6.05
	8	5.60	6.49
	∞	7.54	8.24
Ellipse 	a/b		
	1	3.66	4.36
	2	3.74	4.56
	4	3.79	4.88
	8	3.72	5.09
	16	3.65	5.18
Isosceles Triangle 	θ		
	10°	1.61	2.45
	30°	2.26	2.91
	60°	2.47	3.11
	90°	2.34	2.98
	120°	2.00	2.68

Developing Laminar Flow in the Entrance Region

For a circular tube of length L subjected to constant surface temperature, the average Nusselt number for the *thermal entrance region*:

$$\text{Entry region, laminar:} \quad \text{Nu} = 3.66 + \frac{0.065 (D/L) \text{Re Pr}}{1 + 0.04[(D/L) \text{Re Pr}]^{2/3}}$$

The average Nusselt number is larger at the entrance region, and it approaches asymptotically to the fully developed value of 3.66 as $L \rightarrow \infty$.

When the difference between the surface and the fluid temperatures is large, it may be necessary to account for the variation of viscosity with temperature:

$$\text{Nu} = 1.86 \left(\frac{\text{Re Pr } D}{L} \right)^{1/3} \left(\frac{\mu_b}{\mu_s} \right)^{0.14}$$

All properties are evaluated at the bulk mean fluid temperature, except for μ_s , which is evaluated at the surface temperature.

The average Nusselt number for the thermal entrance region of flow between *isothermal parallel plates* of length L is

$$\text{Entry region, laminar:} \quad \text{Nu} = 7.54 + \frac{0.03 (D_h/L) \text{Re Pr}}{1 + 0.016[(D_h/L) \text{Re Pr}]^{2/3}}$$

$$\text{Re} \leq 2800$$

TURBULENT FLOW IN TUBES

Smooth tubes: $f = (0.790 \ln \text{Re} - 1.64)^{-2}$ $3000 < \text{Re} < 5 \times 10^6$

$\text{Nu} = 0.125 f \text{Re} \text{Pr}^{1/3}$ *Chilton–Colburn analogy*

First Petukhov equation

$$f = 0.184 \text{Re}^{-0.2}$$

$\text{Nu} = 0.023 \text{Re}^{0.8} \text{Pr}^{1/3}$ $\left(\begin{array}{l} 0.7 \leq \text{Pr} \leq 160 \\ \text{Re} > 10,000 \end{array} \right)$ *Colburn equation*

$\text{Nu} = 0.023 \text{Re}^{0.8} \text{Pr}^n$ *Dittus–Boelter equation*

$n = 0.4$ for *heating* and 0.3 for *cooling*

When the variation in properties is large due to a large temperature difference

$$\text{Nu} = 0.027 \text{Re}^{0.8} \text{Pr}^{1/3} \left(\frac{\mu}{\mu_s} \right)^{0.14} \left(\begin{array}{l} 0.7 \leq \text{Pr} \leq 17,600 \\ \text{Re} \geq 10,000 \end{array} \right)$$

All properties are evaluated at T_b except μ_s , which is evaluated at T_s .

$$\text{Nu} = \frac{(f/8) \text{Re} \text{Pr}}{1.07 + 12.7(f/8)^{0.5} (\text{Pr}^{2/3} - 1)} \quad \left(\begin{array}{l} 0.5 \leq \text{Pr} \leq 2000 \\ 10^4 < \text{Re} < 5 \times 10^6 \end{array} \right) \quad \begin{array}{l} \text{Second} \\ \text{Petukhov} \\ \text{equation} \end{array}$$

$$\text{Nu} = \frac{(f/8)(\text{Re} - 1000) \text{Pr}}{1 + 12.7(f/8)^{0.5} (\text{Pr}^{2/3} - 1)} \quad \left(\begin{array}{l} 0.5 \leq \text{Pr} \leq 2000 \\ 3 \times 10^3 < \text{Re} < 5 \times 10^6 \end{array} \right) \quad \begin{array}{l} \text{Gnielinski} \\ \text{relation} \end{array}$$

$$\text{Liquid metals, } T_s = \text{constant:} \quad \text{Nu} = 4.8 + 0.0156 \text{Re}^{0.85} \text{Pr}_s^{0.93}$$

$$\text{Liquid metals, } \dot{q}_s = \text{constant:} \quad \text{Nu} = 6.3 + 0.0167 \text{Re}^{0.85} \text{Pr}_s^{0.93}$$

$$(0.004 < \text{Pr} < 0.01) \quad 10^4 < \text{Re} < 10^6$$

In turbulent flow, wall roughness increases the heat transfer coefficient h by a factor of 2 or more. The convection heat transfer coefficient for rough tubes can be calculated approximately from *Gnielinski relation* or *Chilton–Colburn analogy* by using the friction factor determined from the *Moody chart* or the *Colebrook equation*.

The relations above are not very sensitive to the *thermal conditions* at the tube surfaces and can be used for both $T_s = \text{constant}$ and $q_s = \text{constant}$.

The Moody Chart

The friction factor in fully developed turbulent pipe flow depends on the Reynolds number and the **relative roughness** ε/D .

$$\frac{1}{\sqrt{f}} = -2.0 \log \left(\frac{\varepsilon/D}{3.7} + \frac{2.51}{\text{Re} \sqrt{f}} \right) \quad (\text{turbulent flow})$$

Colebrook equation (for smooth and rough pipes)

$$\frac{1}{\sqrt{f}} \cong -1.8 \log \left[\frac{6.9}{\text{Re}} + \left(\frac{\varepsilon/D}{3.7} \right)^{1.11} \right] \quad \text{Explicit Haaland equation}$$

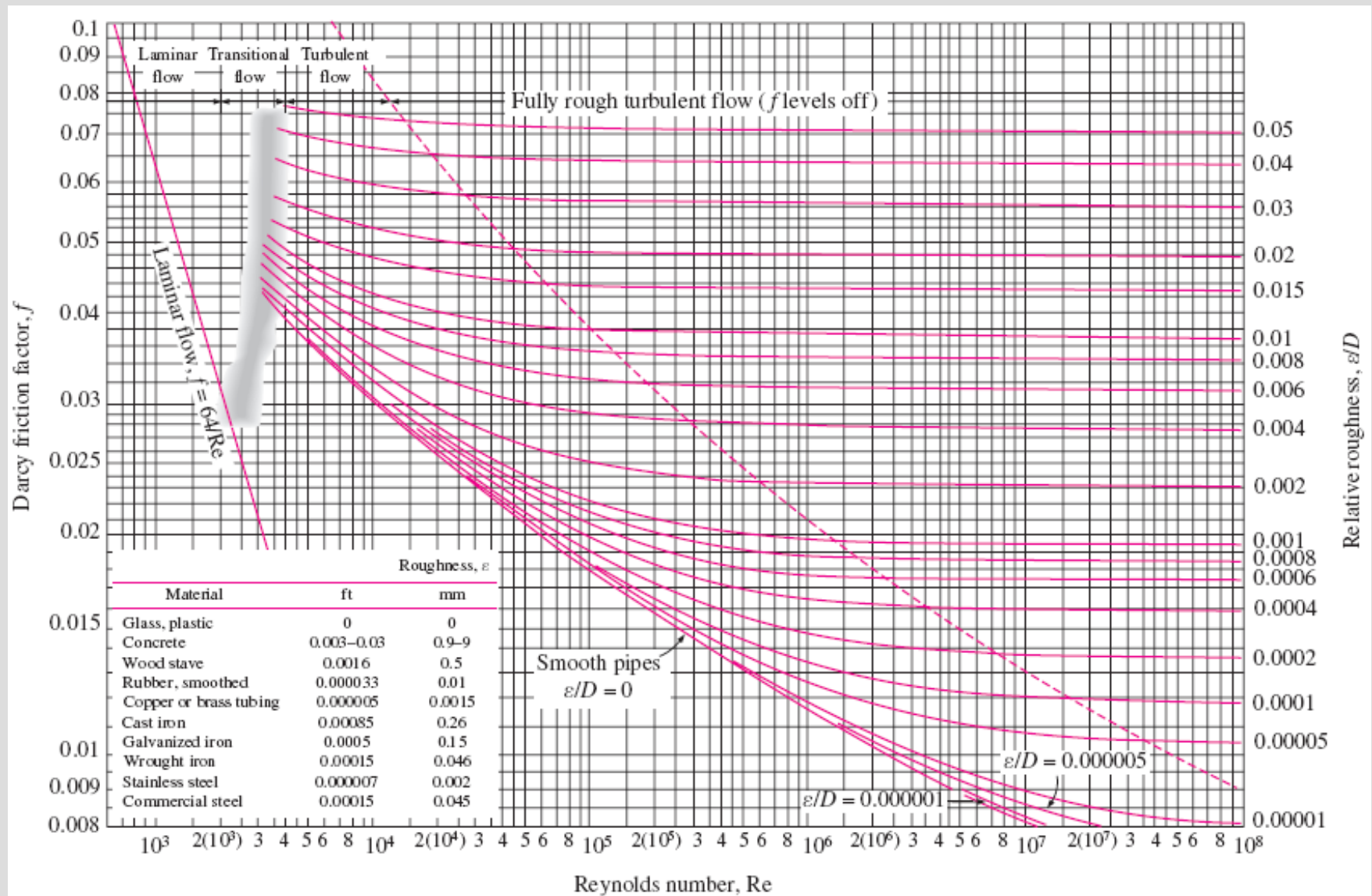
Relative Roughness, ε/D	Friction Factor, f
0.0*	0.0119
0.00001	0.0119
0.0001	0.0134
0.0005	0.0172
0.001	0.0199
0.005	0.0305
0.01	0.0380
0.05	0.0716

* Smooth surface. All values are for $\text{Re} = 10^6$ and are calculated from the Colebrook equation.

The friction factor is minimum for a smooth pipe and increases with roughness.

Equivalent roughness values for new commercial pipes*

Material	Roughness, ε	
	ft	mm
Glass, plastic	0 (smooth)	
Concrete	0.003–0.03	0.9–9
Wood stave	0.0016	0.5
Rubber, smoothed	0.000033	0.01
Copper or brass tubing	0.000005	0.0015
Cast iron	0.00085	0.26
Galvanized iron	0.0005	0.15
Wrought iron	0.00015	0.046
Stainless steel	0.000007	0.002
Commercial steel	0.00015	0.045



The Moody Chart

Developing Turbulent Flow in the Entrance Region

The entry lengths for turbulent flow are typically short, often just 10 tube diameters long, and thus the Nusselt number determined for fully developed turbulent flow can be used approximately for the entire tube.

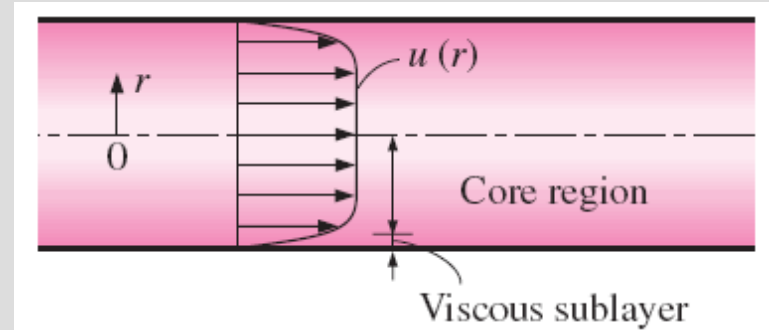
This simple approach gives reasonable results for pressure drop and heat transfer for long tubes and conservative results for short ones.

Correlations for the friction and heat transfer coefficients for the entrance regions are available in the literature for better accuracy.

Turbulent Flow in Noncircular Tubes

Pressure drop and heat transfer characteristics of turbulent flow in tubes are dominated by the very thin viscous sublayer next to the wall surface, and the shape of the core region is not of much significance.

The turbulent flow relations given above for circular tubes can also be used for noncircular tubes with reasonable accuracy by replacing the diameter D in the evaluation of the Reynolds number by the hydraulic diameter $D_h = 4A_c/p$.



In turbulent flow, the velocity profile is nearly a straight line in the core region, and any significant velocity gradients occur in the viscous sublayer. 27

Flow through Tube Annulus

$$D_h = \frac{4A_c}{p} = \frac{4\pi(D_o^2 - D_i^2)/4}{\pi(D_o + D_i)} = D_o - D_i$$

The hydraulic diameter of annulus

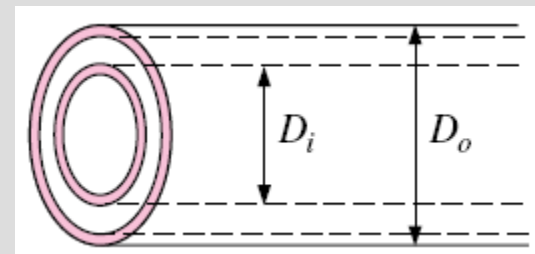
For laminar flow, the convection coefficients for the inner and the outer surfaces are determined from

$$\text{Nu}_i = \frac{h_i D_h}{k} \quad \text{and} \quad \text{Nu}_o = \frac{h_o D_h}{k}$$

For fully developed turbulent flow, h_i and h_o are approximately equal to each other, and the tube annulus can be treated as a noncircular duct with a hydraulic diameter of $D_h = D_o - D_i$. The Nusselt number can be determined from a suitable turbulent flow relation such as the Gnielinski equation. To improve the accuracy, Nusselt number can be multiplied by the following correction factors when one of the tube walls is adiabatic and heat transfer is through the other wall:

$$F_i = 0.86 \left(\frac{D_i}{D_o} \right)^{-0.16} \quad (\text{outer wall adiabatic})$$

$$F_o = 0.86 \left(\frac{D_i}{D_o} \right)^{-0.16} \quad (\text{inner wall adiabatic})$$



Tube surfaces are often *roughened, corrugated, or finned* in order to *enhance* convection heat transfer.

Nusselt number for fully developed laminar flow in an annulus with one surface isothermal and the other adiabatic (Kays and Perkins, 1972)

D_i/D_o	Nu_i	Nu_o
0	—	3.66
0.05	17.46	4.06
0.10	11.56	4.11
0.25	7.37	4.23
0.50	5.74	4.43
1.00	4.86	4.86

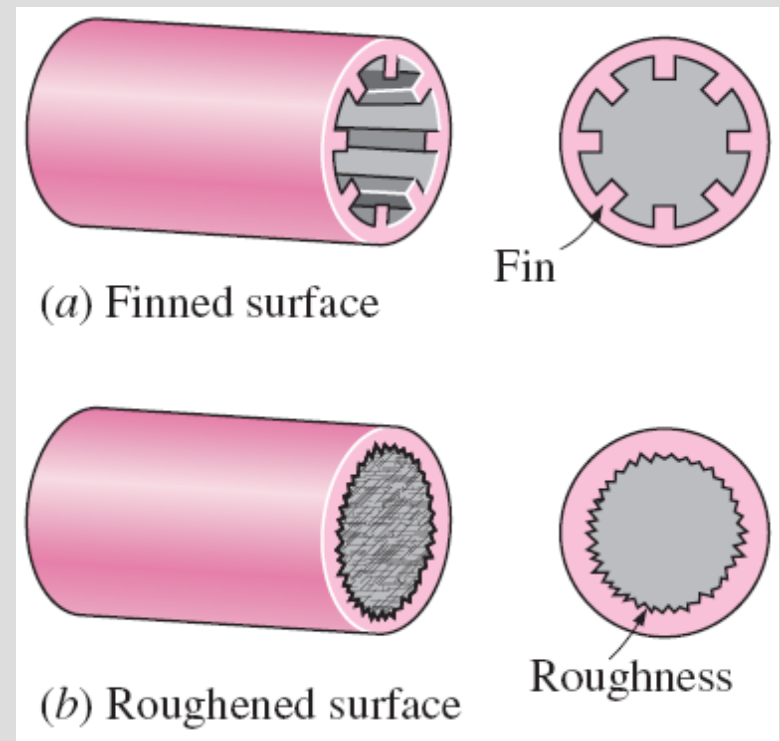
Heat Transfer Enhancement

Tubes with rough surfaces have much higher heat transfer coefficients than tubes with smooth surfaces.

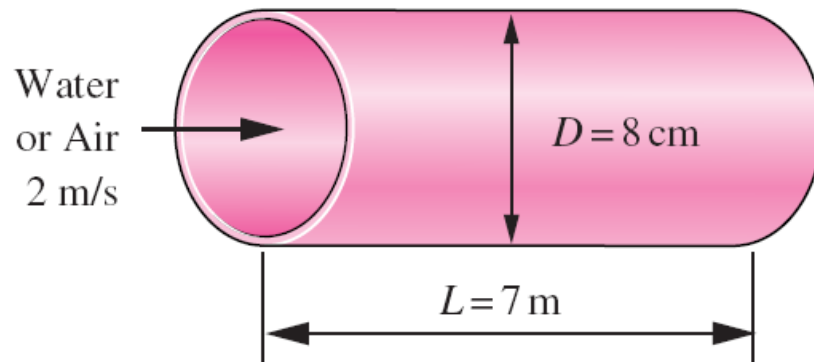
Heat transfer in turbulent flow in a tube has been increased by as much as 400 percent by roughening the surface. Roughening the surface, of course, also increases the friction factor and thus the power requirement for the pump or the fan.

The convection heat transfer coefficient can also be increased by inducing pulsating flow by pulse generators, by inducing swirl by inserting a twisted tape into the tube, or by inducing secondary flows by coiling the tube.

Tube surfaces are often *roughened*, *corrugated*, or *finned* in order to *enhance* convection heat transfer.



8–39 Determine the convection heat transfer coefficient for the flow of (a) air and (b) water at a velocity of 2 m/s in an 8-cm-diameter and 7-m-long tube when the tube is subjected to uniform heat flux from all surfaces. Use fluid properties at 25°C.



Properties The properties of air at 25°C are (Table A-15)

$$k = 0.02551 \text{ W/m}\cdot^\circ\text{C}$$

$$\nu = 1.562 \times 10^{-5} \text{ m}^2/\text{s}$$

$$\text{Pr} = 0.7296$$

The properties of water at 25°C are (Table A-9)

$$\rho = 997 \text{ kg/m}^3$$

$$k = 0.607 \text{ W/m}\cdot^\circ\text{C}$$

$$\nu = \mu / \rho = 0.891 \times 10^{-3} / 997 = 8.937 \times 10^{-7} \text{ m}^2/\text{s}$$

$$\text{Pr} = 6.14$$

Analysis The Reynolds number is

$$\text{Re} = \frac{VD}{\nu} = \frac{(2 \text{ m/s})(0.08 \text{ m})}{1.562 \times 10^{-5} \text{ m}^2/\text{s}} = 10,243$$

which is greater than 10,000. Therefore, the flow is turbulent and the entry lengths in this case are roughly

$$L_h \approx L_t \approx 10D = 10(0.08 \text{ m}) = 0.8 \text{ m}$$

which is much shorter than the total length of the tube. Therefore, we can assume fully developed turbulent flow in the entire duct, and determine the Nusselt number from

$$\text{Nu} = \frac{hD}{k} = 0.023 \text{Re}^{0.8} \text{Pr}^{0.4} = 0.023(10,243)^{0.8} (0.7296)^{0.4} = 32.76$$

Heat transfer coefficient is

$$h = \frac{k}{D} \text{Nu} = \frac{0.02551 \text{ W/m}\cdot^\circ\text{C}}{0.08 \text{ m}} (32.76) = \mathbf{10.45 \text{ W/m}^2\cdot^\circ\text{C}}$$

Repeating calculations for water:

$$\text{Re} = \frac{VD}{\nu} = \frac{(2 \text{ m/s})(0.08 \text{ m})}{8.937 \times 10^{-7} \text{ m}^2/\text{s}} = 179,035$$

$$\text{Nu} = \frac{hD}{k} = 0.023 \text{Re}^{0.8} \text{Pr}^{0.4} = 0.023(179,035)^{0.8} (6.14)^{0.4} = 757.4$$

$$h = \frac{k}{D} \text{Nu} = \frac{0.607 \text{ W/m}\cdot^\circ\text{C}}{0.08 \text{ m}} (757.4) = \mathbf{5747 \text{ W/m}^2\cdot^\circ\text{C}}$$

Discussion The heat transfer coefficient for water is 550 times that of air.

8–88 Cold air at 5°C enters a 12-cm-diameter 20-m-long isothermal pipe at a velocity of 2.5 m/s and leaves at 19°C. Estimate the surface temperature of the pipe.

Properties The properties of air at 1 atm and the bulk mean temperature of $(5+19)/2=12^\circ\text{C}$ are (Table A-15)

$$\rho = 1.238 \text{ kg/m}^3$$

$$k = 0.02454 \text{ W/m}\cdot^\circ\text{C}$$

$$\nu = 1.444 \times 10^{-5} \text{ m}^2/\text{s}$$

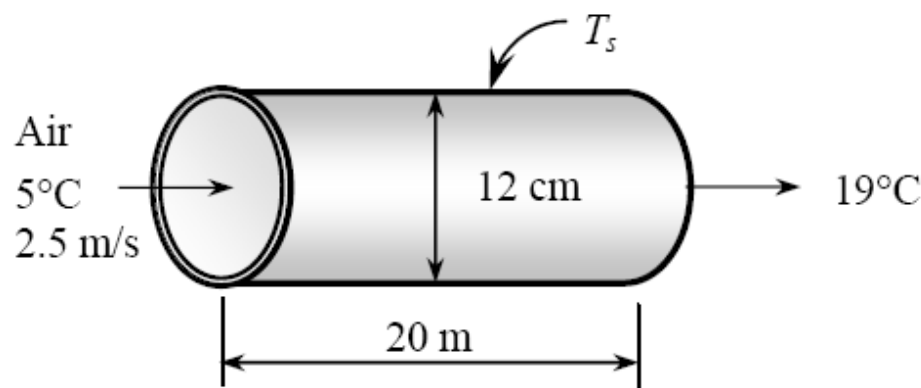
$$c_p = 1007 \text{ J/kg}\cdot^\circ\text{C}$$

$$\text{Pr} = 0.7331$$

Analysis The rate of heat transfer to the air is

$$\dot{m} = \rho A_c V_{\text{avg}} = (1.238 \text{ kg/m}^3) \pi \frac{(0.12 \text{ m})^2}{4} (2.5 \text{ m/s}) = 0.0350 \text{ kg/s}$$

$$\dot{Q} = \dot{m} c_p \Delta T = (0.0350 \text{ kg/s})(1007 \text{ J/kg}\cdot^\circ\text{C})(19 - 5)^\circ\text{C} = 493.1 \text{ W}$$



Reynolds number is

$$\text{Re} = \frac{V_{\text{avg}} D}{\nu} = \frac{(2.5 \text{ m/s})(0.12 \text{ m})}{1.444 \times 10^{-5} \text{ m}^2/\text{s}} = 20,775$$

which is greater than 10,000. Therefore, the flow is turbulent and the entry lengths in this case are roughly

$$L_h \approx L_t \approx 10D = 10(0.12 \text{ m}) = 1.2 \text{ m}$$

which is much shorter than the total length of the duct. Therefore, we can assume fully developed turbulent flow in the entire duct, and determine the Nusselt number from

$$\text{Nu} = \frac{hD}{k} = 0.023 \text{Re}^{0.8} \text{Pr}^{0.4} = 0.023(20,775)^{0.8} (0.7331)^{0.4} = 57.79$$

Heat transfer coefficient is

$$h = \frac{k}{D} \text{Nu} = \frac{0.02454 \text{ W/m} \cdot ^\circ\text{C}}{0.12 \text{ m}} (57.79) = 11.82 \text{ W/m}^2 \cdot ^\circ\text{C}$$

The logarithmic mean temperature difference is determined from

$$\dot{Q} = hA_s \Delta T_{\text{ln}} \longrightarrow 493.1 \text{ W} = (11.82 \text{ W/m}^2 \cdot ^\circ\text{C}) [\pi(0.12 \text{ m})(20 \text{ m})] \Delta T_{\text{ln}} \longrightarrow \Delta T_{\text{ln}} = 5.533^\circ\text{C}$$

Then the pipe temperature is determined from the definition of the logarithmic mean temperature difference

$$\Delta T_{\text{ln}} = \frac{T_e - T_i}{\ln\left(\frac{T_s - T_e}{T_s - T_i}\right)} \longrightarrow 5.533^\circ\text{C} = \frac{19 - 5}{\ln\left(\frac{T_s - 19}{T_s - 5}\right)} \longrightarrow T_s = 3.8^\circ\text{C}$$

Summary

- General Considerations for Pipe Flow
 - ✓ Thermal Entrance Region, Entry Lengths
- General Thermal Analysis
 - ✓ Constant Surface Heat Flux
 - ✓ Constant Surface Temperature
- Laminar Flow in Tubes
 - ✓ Constant Surface Heat Flux, Constant Surface Temperature
 - ✓ Laminar Flow in Noncircular Tubes, Developing Laminar Flow in the Entrance Region
- Turbulent Flow in Tubes
 - ✓ Developing Turbulent Flow in the Entrance Region,
 - ✓ Turbulent Flow in Noncircular Tubes
 - ✓ Flow through Tube Annulus, Heat Transfer Enhancement